

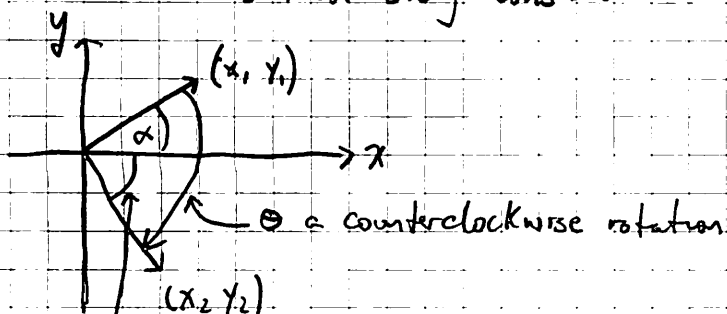
# Solutions to PS # 2

1) a)  $\begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} = E$  This matrix was explicitly derived in class

b) For proper rotation around any  $C_n$  Axis, the following analysis will hold:

\* By convention, the  $C_n$  axis is the  $z$ -axis, so our  $z$  coordinate will not change upon rotation

\* We need only consider how  $x + y$  project upon rotation



$-(\theta - \alpha) \Rightarrow$  This angle is negative because it is measured from the  $x$  axis in the opposite direction of  $\alpha$

$$\begin{aligned} x_1 &= \vec{v} \cos \alpha & \xrightarrow{C_n(\theta)} & x_2 = \vec{v} \cos [-(\theta - \alpha)] = -\vec{v} \cos(\theta - \alpha) \\ y_1 &= \vec{v} \sin \alpha & & y_2 = \vec{v} \sin [-(\theta - \alpha)] = -\vec{v} \sin(\theta - \alpha) \end{aligned}$$

Using the identity operation we get the following expression...

$$x_2 = -\vec{v} \cos(\theta - \alpha) = -\vec{v} [\cos \theta \cos \alpha + \sin \theta \sin \alpha] = -x_1 \cos \theta - y_1 \sin \theta$$

$$y_2 = -\vec{v} \sin(\theta - \alpha) = -\vec{v} [\sin \theta \cos \alpha - \cos \theta \sin \alpha] = -x_1 \sin \theta + y_1 \cos \theta$$

If we reformulate the above expressions in terms of a matrix representation, we get:

$$C_n(\theta) \begin{bmatrix} x_1 \\ y_1 \\ z_1 \end{bmatrix} = \begin{bmatrix} -x_1 \cos \theta - y_1 \sin \theta \\ -x_1 \sin \theta + y_1 \cos \theta \\ z_1 \end{bmatrix} \quad \therefore C_n(\theta) = \begin{bmatrix} -\cos \theta & -\sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad \text{where } \theta = \frac{2\pi}{n}$$

The expression above is general for any  $C_n$  rotation.

For the case of  $C_3$

$$\begin{bmatrix} 1/2 & -\sqrt{3}/2 & 0 \\ -\sqrt{3}/2 & -1/2 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

c) Using the same derivation + analysis above, we find

$$C_y = \begin{bmatrix} 0 & -1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

d)

$$L = \begin{bmatrix} -1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{bmatrix}$$

This matrix was expressly derived in class

e)

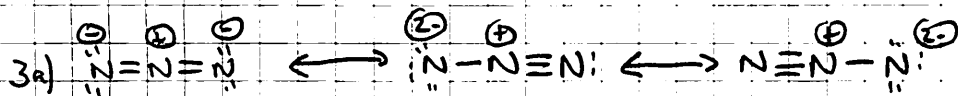
$$O_h \cdot C_n(\theta) = S_n(\theta)$$

$$\begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & -1 \end{bmatrix} \begin{bmatrix} -\cos \theta & \sin \theta & 0 \\ -\sin \theta & -\cos \theta & 0 \\ 0 & 0 & 1 \end{bmatrix} = \begin{bmatrix} -\cos \theta & -\sin \theta & 0 \\ -\sin \theta & \cos \theta & 0 \\ 0 & 0 & -1 \end{bmatrix}$$

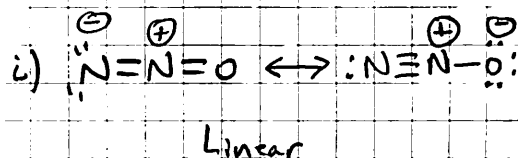
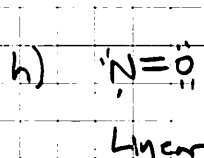
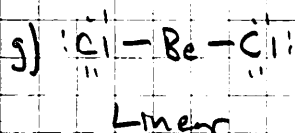
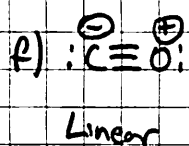
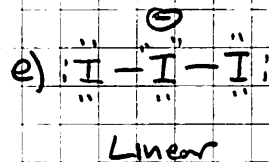
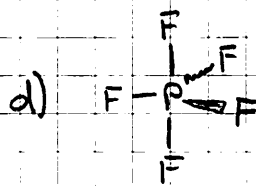
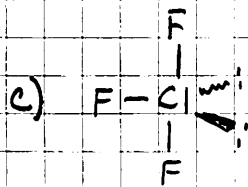
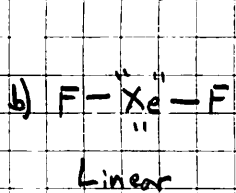
For  $S_4$ ,  $\theta = \frac{2\pi}{4} \therefore S_4 = \begin{bmatrix} 0 & -1 & 0 \\ -1 & 0 & 0 \\ 0 & 0 & -1 \end{bmatrix}$

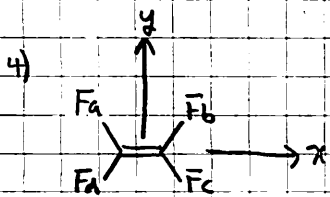
2a)  $C_{2v}$  b)  $C_{3v}$  c)  $C_s$  d)  $D_{2h}$  e)  $D_{2d}$  f)  $O_h$  g)  $D_{6h}$  h)  $T_d$  i)  $T_d$

j)  $O_h$  k)  $D_{2d}$  l)  $D_{2h}$  m)  $C_s$  n)  $D_{3h}$  o)  $D_{2d}$  p)  $C_{2v}$  q)  $D_3$



Linear

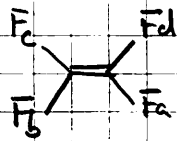




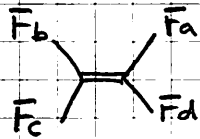
$D_{2h}$	E	$C_2(z)$	$C_2(y)$	$C_2(x)$	<del><math>S_4</math></del>	i	$\sigma(xy)$	$\sigma(xz)$	$\sigma(yz)$
$F_a$	$F_a$	$F_c$	$F_b$	$F_d$	↑	$F_c$	$F_a$	$F_d$	$F_b$

$S_4$  should not be included!

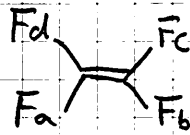
$C_2(z)$



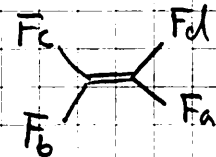
$C_2(y)$



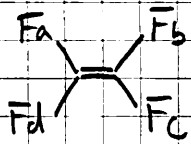
$C_2(x)$



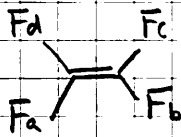
i



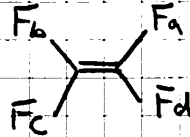
$\sigma(xy)$



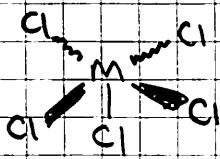
$\sigma(xz)$



$\sigma(yz)$

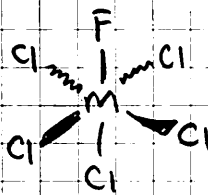


5a)



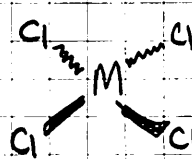
$C_{4v}$

b)



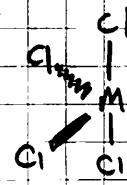
$C_{4v}$

c)



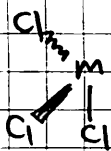
$D_{4h}$

d)



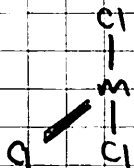
$C_{2v}$

e)



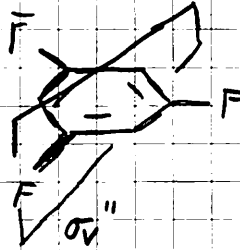
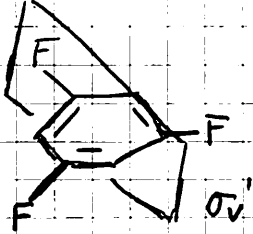
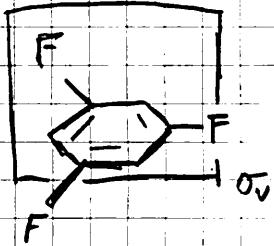
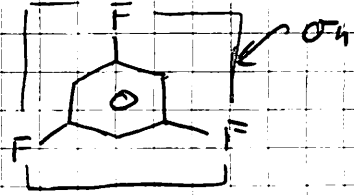
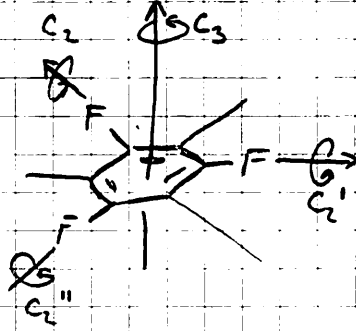
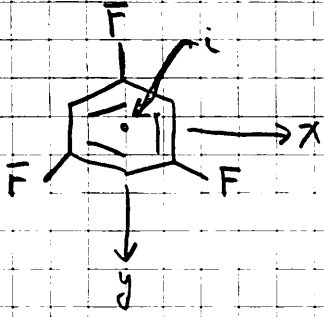
$C_{3v}$

f)



$C_{2v}$

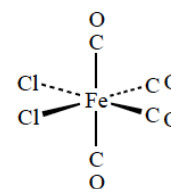
6)



4.28 a. *cis*-Fe(CO)<sub>4</sub>Cl<sub>2</sub> has C<sub>2v</sub> symmetry.

The vectors for CO stretching have the representation  $\Gamma$ :

C <sub>2v</sub>	E	C <sub>2</sub>	$\sigma_v(xz)$	$\sigma_v'(yz)$	
$\Gamma$	4	0	2	2	
A <sub>1</sub>	1	1	1	1	z
A <sub>2</sub>	1	1	-1	-1	
B <sub>1</sub>	1	-1	1	-1	x
B <sub>2</sub>	1	-1	-1	1	y



$$n(A_1) = 1/4[4 \times 1 + 0 \times 1 + 2 \times 1 + 2 \times 1] = 2$$

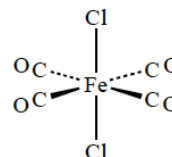
$$n(A_2) = 1/4[4 \times 1 + 0 \times 1 + 2 \times (-1) + 2 \times (-1)] = 0$$

$$n(B_1) = 1/4[4 \times 1 + 0 \times (-1) + 2 \times 1 + 2 \times (-1)] = 1$$

$$n(B_2) = 1/4[4 \times 1 + 0 \times (-1) + 2 \times (-1) + 2 \times 1] = 1$$

$$\Gamma = 2A_1 + B_1 + B_2, \text{ all four IR active.}$$

b. *trans*-Fe(CO)<sub>4</sub>Cl<sub>2</sub> has D<sub>4h</sub> symmetry.



D <sub>4h</sub>	E	2C <sub>4</sub>	C <sub>2</sub>	2C <sub>2</sub> '	2C <sub>2</sub> ''	i	2S <sub>4</sub>	$\sigma_h$	2 $\sigma_v$	2 $\sigma_d$	
$\Gamma$	4	0	0	2	0	0	0	4	2	0	
A <sub>2u</sub>	1	1	1	-1	-1	-1	-1	-1	1	1	z
E <sub>u</sub>	2	0	-2	0	0	-2	0	2	0	0	(x,y)

Omitting the operations that have zeroes in  $\Gamma$ :

$$n(A_{2u}) = 1/16[4 \times 1 + 2 \times 2 \times (-1) + 4 \times (-1) + 2 \times 2 \times 1] = 0$$

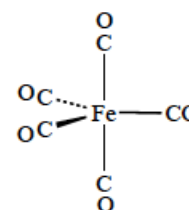
$$n(E_u) = 1/16[4 \times 2 + 2 \times 2 \times 0 + 4 \times 2 + 2 \times 2 \times 0] = 1 \text{ (IR active)}$$

Note: In checking for IR-active bands, it is only necessary to check the irreducible representations having the same symmetry as x, y, or z, or a combination of them.

c. Fe(CO)<sub>5</sub> has D<sub>3h</sub> symmetry.

The vectors for C–O stretching have the following representation  $\Gamma$ :

D <sub>3h</sub>	E	2C <sub>3</sub>	3C <sub>2</sub>	$\sigma_h$	2S <sub>3</sub>	3 $\sigma_v$	
$\Gamma$	5	2	1	3	0	3	
E'	2	-1	0	2	-1	0	(x,y)
A <sub>2</sub> ''	1	1	-1	-1	-1	1	z



$$n(E') = 1/12 [(5 \times 2) + (2 \times 2 \times -1) + (3 \times 2)] = 1$$

$$n(A_2'') = 1/12 [(5 \times 1) + (2 \times 2 \times 1) + (3 \times 1 \times -1) + (3 \times -1) + (3 \times 3 \times 1)] = 1$$

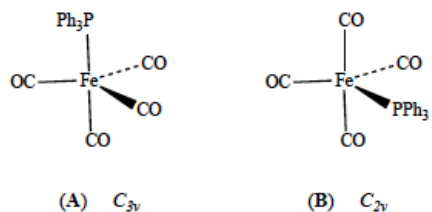
There are two bands, one matching E' and one matching A<sub>2</sub>''. These are the only irreducible representations that match the coordinates x, y, and z.

- 4.29 In 4.28a, the symmetries of the CO stretching vibrations of *cis*-Fe(CO)<sub>4</sub>Cl<sub>2</sub> (*C*<sub>2v</sub> symmetry) are determined as 2 *A*<sub>1</sub> + *B*<sub>1</sub> + *B*<sub>2</sub>. Each of these representations matches Raman-active functions: *A*<sub>1</sub> (*x*<sup>2</sup>, *y*<sup>2</sup>, *z*<sup>2</sup>); *A*<sub>2</sub> (*xy*); *B*<sub>1</sub> (*xz*); and *B*<sub>2</sub> (*yz*), so all are Raman-active.

In 4.28b, the symmetries of the CO stretching vibrations of *trans*-Fe(CO)<sub>4</sub>Cl<sub>2</sub> (*D*<sub>4h</sub> symmetry) are *A*<sub>1g</sub> + *B*<sub>1g</sub> + *E*<sub>u</sub>. Only *A*<sub>1g</sub> (*x*<sup>2</sup> + *y*<sup>2</sup>, *z*<sup>2</sup>) and *B*<sub>1g</sub> (*x*<sup>2</sup> - *y*<sup>2</sup>) match Raman active functions; this complex exhibits two Raman-active CO stretching vibrations.

In 4.28c, the symmetries of the CO stretching vibrations of Fe(CO)<sub>5</sub> (*D*<sub>3h</sub> symmetry) are 2 *A*<sub>1</sub>' + *E*' + *A*<sub>2</sub>'. Only *A*<sub>1</sub>' (*x*<sup>2</sup> + *y*<sup>2</sup>, *z*<sup>2</sup>) and *E*' (*x*<sup>2</sup> - *y*<sup>2</sup>, *xy*) match Raman-active functions; this complex exhibits four Raman-active CO stretching vibrations.

- 4.33 The possible isomers are as follows, with the triphenylphosphine ligand in either the axial (A) or equatorial (B) sites.



Note that the triphenylphosphine ligand is approximated as a simple L ligand for the sake of the point group determination. Rotation about the Fe–P bond in solution is expected to render the arrangement of the phenyl rings unimportant in approximating the symmetry of these isomers in solution. The impact of the phenyl rings would likely be manifest in the IR  $\nu(\text{CO})$  spectra of these isomers in the solid-state.

For A, we consider each CO bond as a vector to deduce the expected number of carbonyl stretching modes. The irreducible representation is as follows:

<i>C</i> <sub>3v</sub>	<i>E</i>	<i>C</i> <sub>3</sub>	3 $\sigma_v$		
$\Gamma$	4	1	2		
<i>A</i> <sub>1</sub>	1	1	1	<i>z</i>	<i>x</i> <sup>2</sup> + <i>y</i> <sup>2</sup> , <i>z</i> <sup>2</sup>
<i>A</i> <sub>2</sub>	1	1	-1	<i>R</i> <sub>x</sub>	
<i>E</i>	2	-1	0	( <i>x</i> , <i>y</i> ), ( <i>R</i> <sub>x</sub> , <i>R</i> <sub>y</sub> )	( <i>x</i> <sup>2</sup> - <i>y</i> <sup>2</sup> , <i>xy</i> ), ( <i>xz</i> , <i>yz</i> )

Reduction of the reducible representation affords 2 *A*<sub>1</sub> + *E*. These stretching modes are IR-active and three  $\nu(\text{CO})$  absorptions are expected for A.

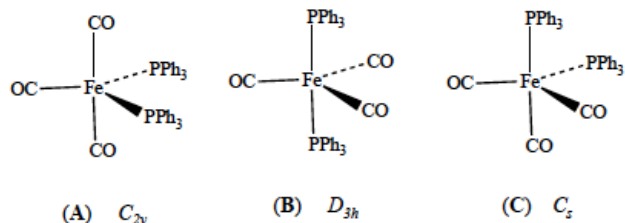
For B, a similar analysis affords the following irreducible representation:

<i>C</i> <sub>2v</sub>	<i>E</i>	<i>C</i> <sub>2</sub>	$\sigma_v(xz)$	$\sigma_v'(yz)$		
$\Gamma$	4	0	2	2		
<i>A</i> <sub>1</sub>	1	1	1	1	<i>z</i>	<i>x</i> <sup>2</sup> , <i>y</i> <sup>2</sup> , <i>z</i> <sup>2</sup>
<i>A</i> <sub>2</sub>	1	1	-1	-1	<i>R</i> <sub>z</sub>	<i>xy</i>
<i>B</i> <sub>1</sub>	1	-1	1	-1	<i>x</i> , <i>R</i> <sub>y</sub>	<i>xz</i>
<i>B</i> <sub>2</sub>	1	-1	-1	1	<i>y</i> , <i>R</i> <sub>x</sub>	<i>yz</i>

Reduction of the reducible representation affords 2 *A*<sub>1</sub> + *B*<sub>1</sub> + *B*<sub>2</sub>. These stretching modes are IR-active, and four  $\nu(\text{CO})$  absorptions are expected for A.

The reported  $\nu(\text{CO})$  IR spectrum is consistent with formation of isomer A, with the triphenylphosphine ligand in the axial site.

4.34 As in 4.33, we consider the triphenylphosphine ligand as a simple L group for point group determination. The point groups for isomers A, B, and C are as follows:



For A, the set of irreducible representations for the three CO stretching vibrational modes is  $2A_1 + B_1$ . These modes are all IR-active in the  $C_{2v}$  character table, and three  $\nu(\text{CO})$  IR absorptions are expected for isomer A.

For B, the set of irreducible representations for the three CO stretching vibrational modes is  $A_1' + E'$ . Only the  $E'$  mode is IR-active in the  $D_{3h}$  point group, and one  $\nu(\text{CO})$  IR absorption is expected for isomer B.

For C, the set of irreducible representations for the three CO stretching vibrational modes is  $2A' + A''$ . These modes are all IR-active in the  $C_s$  point group, and three  $\nu(\text{CO})$  IR absorptions are expected for isomer C.

The single  $\nu(\text{CO})$  IR absorption reported for  $\text{Fe}(\text{CO})_3(\text{PPh}_3)_2$  supports the presence of the  $D_{3h}$  isomer B.

The *trans* isomer B is reported in R. L. Keiter, E. A. Keiter, K. H. Hecker, C. A. Boecker, *Organometallics*, 1988, 7, 2466, and the authors observe splitting of the absorption associated with the  $E'$  mode in  $\text{CHCl}_3$ . The forbidden  $A_1'$  stretching mode was observed as a weak absorption.